

*Universitat Politecnica De Catalunya, BarcelonaTech  
Masters in Computational Mechanics*

*Course  
Computational Structural Mechanics and Dynamics*

**Assignment 9**  
**on**  
**Axisymmetric Shells**

by

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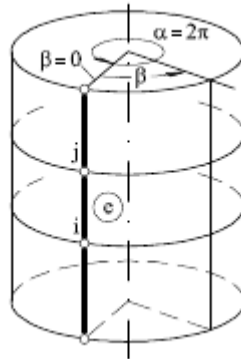
**1. Describe in extension how can be applied a non-symmetric load on this formulation?**

The axisymmetric shell formulation can be extended for non-symmetric loading. The displacement field is split into symmetric and anti-symmetric component with respect to plane at  $\beta=0$ . For a n noded strip with n nodes, we have,

$$u' = \sum_{l=0}^m \sum_{i=1}^n N_i (\bar{S}^l \bar{a}'_i + \bar{S}^l \bar{a}'_i)$$

where  $u'$  is the displacement vector,  $(\bar{\cdot})$  and  $(\bar{\cdot})$  axisymmetric and anti-symmetric components of displacements.

$$\bar{S}^l(y) = \begin{bmatrix} S^l & & & 0 \\ & C^l & & \\ & & S^l & \\ 0 & & & C^l \end{bmatrix}$$



$S^l = \sin(\gamma y)$  where  $\gamma = l/b$ .  $b$  is the plate length  $\bar{S}^l$  is where  $S^l$  is replaced by  $C^l$  and visa-versa.

The loads are expanded in Fourier series using harmonic functions as for the displacements.

$$t = \sum_{l=0}^m (\bar{S}^l \bar{t}' + \bar{S}^l \bar{t}'')$$

The local stiffness matrix for axisymmetric strip element is given by

$$[K'_{ij}] = C \int_{d^{(e)}} [B'_i]^T \hat{D}' B'_j r ds$$

where,

$$C = \begin{cases} 2\pi & \text{for } l=0 \\ \pi & \text{for } l \neq 0 \end{cases}$$

For symmetric case,  $\gamma$  is replaced by  $-l$ .

For anti-symmetric case,  $\gamma$  is replaced by  $l$ .

In global stiffness matrix is,

$$[K^{ll}_{ij}] = C \int_{d^{(e)}} [B_i^l]^T \hat{D}' B_j^l r ds \quad \text{where } B_i^l = B'_i L_i^{(e)} \quad L_i^{(e)} \text{ is the transformation matrix.}$$

**2. Using thin beams formulation, describe the shape of the B(e) matrix and comment the integration rule.**

Using thin beam formulation leads to neglecting the effect of transverse shear strains in the analysis. The formulation is therefore applicable to thin shell problems only. Making  $\hat{\epsilon}_s=0$  we get,

$$\hat{\epsilon}' = \sum_{l=1}^m \sum_{i=1}^n \hat{S}^l B_i'^l a_i'^l ; \quad a_i'^l = [u_0'^l, v_0'^l, w_0'^l, \frac{\partial w_0'^l}{\partial s}]$$

$$\hat{\epsilon}' = \begin{pmatrix} \hat{\epsilon}'_m \\ \hat{\epsilon}'_b \end{pmatrix} ; \quad \hat{\epsilon}'_m = \begin{pmatrix} \frac{\partial u_0'}{\partial s} \\ \frac{1}{r} \frac{\partial v_0'}{\partial \beta} + \frac{u_0'}{r} C - \frac{w_0'}{r} \\ \frac{\partial v_0'}{\partial s} + \frac{1}{r} \frac{\partial u_0'}{\partial \beta} - \frac{v_0'}{r} C \end{pmatrix} ; \quad \hat{\epsilon}'_b = \begin{pmatrix} \frac{\partial^2 w_0'}{\partial s^2} \\ \frac{1}{r} \frac{\partial^2 w_0'}{\partial \beta^2} + \frac{S}{r^2} \frac{\partial v_0'}{\partial \beta} + \frac{C}{r} \frac{\partial w_0'}{\partial s} \\ \frac{2}{r} \frac{\partial^2 w_0'}{\partial s \partial \beta} - \frac{2CS}{r^2} v_0' - \frac{2C}{r^2} \frac{\partial w_0'}{\partial \beta} + \frac{S}{r} \frac{\partial v_0'}{\partial s} \end{pmatrix}$$

$$B'^l_{m_i} = \begin{pmatrix} B'^l_{m_i} \\ B'^l_{b_i} \end{pmatrix} ; \quad B'^l_{m_i} = \begin{bmatrix} \frac{\partial N_i}{\partial s} & 0 & 0 & 0 \\ \frac{N_i}{r} C & -\frac{N_i}{r} \gamma & -\frac{H_i'}{r} S & -\frac{\bar{H}_i}{r} S \dot{\gamma} \\ \frac{N_i}{r} C & (\frac{\partial N_i}{\partial s} - \frac{N_i}{r} C) & 0 & 0 \end{bmatrix} ;$$

$$B'^l_{b_i} = \begin{bmatrix} 0 & 0 & \frac{\partial^2 H_i}{\partial s^2} & \frac{\partial^2 \bar{H}_i}{\partial s^2} \\ 0 & \frac{N_i}{r^2} S \gamma & [\frac{C}{r} \frac{\partial H_i}{\partial s} s - (\frac{\gamma}{r})^2 H_i] & [\frac{C}{r} \frac{\partial \bar{H}_i}{\partial s} s - (\frac{\gamma}{r})^2 \bar{H}_i] \\ 0 & (\frac{S}{r} \frac{\partial N_i}{\partial s} - 2 \frac{N_i}{r^2} C S) & (2 \frac{\gamma}{r} \frac{\partial H_i}{\partial s} - 2 \frac{H_i}{r^2} C \gamma) & (2 \frac{\gamma}{r} \frac{\partial \bar{H}_i}{\partial s} - 2 \frac{\bar{H}_i}{r^2} C \gamma) \end{bmatrix}$$

$N_i$  = Langrange Shape Functions ;  $H_i, \bar{H}_i$  = Hermite Shape Functions

**Integration Rule:** Two-point quadrature is recommended for computing the integrals. But results also can be obtained using simplest reduced one-point quadrature.